The Fourier Pseudospectral Method with a Restrain Operator for the RLW Equation

BEN-YU GUO AND WEI-MING CAO

Shanghai University of Science and Technology, Shanghai, China

Received December 12, 1985; Revised March 13, 1987

In this paper we develop a new Fourier pseudospectral method with a restrain operator which is applied to the RLW equation. The numerical results show the advantages of this method. We prove the generalized stability and the convergence of the scheme. © 1988 Academic Press, Inc.

I. INTRODUCTION

In this paper we consider the periodical problem of the RLW equation

$$\frac{\partial U}{\partial t} + \alpha \frac{\partial U}{\partial x} + U \frac{\partial U}{\partial x} - \delta \frac{\partial^3 U}{\partial t \partial x^2} = f, \qquad x \in \vartheta; \ t \in (0, T],$$

$$\frac{\partial^q U(0, t)}{\partial x^q} = \frac{\partial^q U(2, t)}{\partial x^q}, \qquad q = 0, 1; \ t \in [0, T], \qquad (1)$$

$$U(x, 0) = U_0(x), \qquad x \in \vartheta,$$

where $\vartheta = (0, 2), \ \delta > 0, \ f(0, t) = f(2, t)$, and α is a constant.

In 1966, Peregrine [1] proposed the first numerical method for solving (1). In 1976, Abdulloev, Bogolubsky, Makhankov [2] showed numerically that the collision between two soliton-like waves is inelastic. Olver [3] proved that a RLW equation possesses only three conservations.

Other numerical methods can be found in Eilbeck and McGuire [4,5], Alexander and Morris [6], and Arnold *et al.*, [7]. Recently Wu and Guo [8], using a difference scheme with high accuracy, found that four waves were formed after the collision between two soliton-like waves. The accuracy of numerical solution of both finite element and difference methods is limited, even if the solution of (1) is very smooth.

The spectral method and pseudospectral method are two important numerical methods for P.D.E. (see Gottlieb and Orszag [9]). Pasciak [10] and Guo and Manoranjan [11] used these methods to solve the RLW equation. The pseudo-spectral method is preferable when dealing with nonlinear terms. However, it is less stable than the spectral method due to the aliasing interaction. For remedying this

deficiency Kreiss and Oliger [12], Kuo [13], and Ma and Guo Ben-yu [14] proposed several filtering techniques. In this paper, we modify the filtering technique used in [14] and give a new pseudospectral method applied to (1). We test numerically the effect of the restrain operator and find that

(i) If the vibration of the solution of (1) is small and δ large, then the effect of the restrain operator is not clear.

(ii) If the vibration of the solution of (1) is big or δ very small, then the restrain operator increases the stability.

(iii) If we use the filtering technique in the same way as in [14] (i.e., scheme (8) in this paper), then the numerical result is poor, especially for very small δ (see Table VI, in Section III). But the new method proposed here gives a uniformly good approximation independent of δ (see Tables V, VI, in Section III).

We use the technique of generalized stability of Guo [15, 16] (i.e., g-stability; see Griffiths [17]), to strictly estimate the error. Four estimates are established. Only two of them depend on δ . The convergence is given also.

II. THE SCHEME

Let

$$(v, w) = \int_0^2 v(x) \overline{w(x)} dx,$$
$$\|v\|^2 = (v, v), \qquad |v|_1 = \left\|\frac{\partial v}{\partial x}\right\|$$

Now we derive the new pseudospectral method. Let N be a positive integer and n any integer:

$$V_n = \operatorname{Span} \{ e^{n\pi i x} / |n| \leq N \}.$$

 \dot{V}_N is a real value function subset of V_N . Let P_N be the orthogonal projection operator from $L^2(\vartheta)$ to V_N , i.e.,

$$(p_N v, w) = (v, w), \quad \forall w \in V_N.$$

Let P_c be the interpolation operator from $C(\mathfrak{G})$ to V_N such that

$$p_c v(x_j) = v(x_j), \qquad x_j = \frac{2j}{2N+1}, \qquad j = 0, 1, ..., 2N.$$

From [12],

$$(v, w)_N = (p_c v, p_c w)_N = (p_c v, p_c w), \qquad \forall v, w \in c(\mathfrak{S}),$$
(2)

where

$$(v, w)_N = \frac{2}{2N+1} \sum_{j=0}^{2N} v(x_j) \overline{w(x_j)}.$$

Let τ be the mesh spacing of the variable t, $S_{\tau} = \{k \mid 0 \le k \le \lfloor T/\tau \rfloor - 1\}$ and $v^{k}(x) = v(x, k\tau)$. Define

$$v_t^k(x) = \frac{1}{\tau} (v^{k+1}(x) - v^k(x))$$

As is well known, a reasonable scheme must keep the properties similar to those of (1). Indeed the solution of (1) satisfies the following conservations

$$\int_0^2 U(x,t) \, dx = \int_0^2 U_0(x) \, dx + \int_0^t \int_0^2 f(x,\xi) \, dx \, d\xi. \tag{3}$$

$$\|U(t)\|^{2} + \delta \|U(t)\|_{1}^{2} = \|U_{0}\|^{2} + \delta \|U_{0}\|_{1}^{2} + 2\int_{0}^{t} (f(\xi), U(\xi)) d\xi.$$
(4)

In order to simulate (3) and (4), we define the following operator

$$J(v(x), \varphi(x)) = \frac{1}{3} p_c \left(\varphi(x) \frac{\partial}{\partial x} v(x)\right) + \frac{1}{3} \frac{\partial}{\partial x} p_c(\varphi(x), v(x)).$$

It can be shown that for all $v, w, \varphi \in \dot{V}_N$,

$$(J(v,\varphi),\omega) + (J(w,\varphi),v) = 0.$$
⁽⁵⁾

The simplest pseudospectral method for (1) is to find $u^k \in \dot{V}_N$ such that

$$u_{t}^{k}(x) + \alpha \frac{\partial}{\partial x} u^{k}(x) + J(u^{k}(x) + \sigma \tau u_{t}^{k}(x), u^{k}(x)) - \delta \frac{\partial^{2}}{\partial x^{2}} u_{t}^{k}(x)$$

$$= p_{c} f^{k}(x), \qquad x \in \vartheta, k \in s_{\tau},$$

$$u^{0}(x) = P_{c} U_{0}(x), \qquad x \in \vartheta,$$
(6)

where $0 \le \delta \le 1$. But the numerical results showed that scheme (6) is not stable for very small δ . On the other hand, for the error estimation uniformly for δ , we have to establish the following inequality

$$\left| \left(F\left(v \frac{\partial v}{\partial x}\right), w \right)_{N} \right| \leq b(w) \|v\|^{2}, \quad \forall v, w \in V_{N},$$
(7)

where b(w) is independent of N, and

$$F = \left(I - \delta \frac{\partial^2}{\partial x^2}\right)^{-1}.$$

112

But the inequality (7) is not true generally. For instance, we consider the following functions

$$v(x) = a_N e^{N\pi i x}, \qquad w(x) = 1 + 2i \sin \pi x,$$

and so from (2)

$$\left(F\left(v\frac{\partial v}{\partial x}\right), w\right)_{N} = \left(p_{c}F\left(v\frac{\partial v}{\partial x}\right), w\right).$$

Since

$$\begin{split} \left| \left(p_c F\left(v \frac{\partial v}{\partial x} \right), w \right) \right| &= \frac{\pi N a_N^2}{1 + \delta \pi^2 N^2} \left| \left(P_c e^{2N\pi i x}, w \right) \right| \\ &= \frac{\pi N a_N^2}{1 + \delta \pi^2 N^2} \left| \left(P_c e^{-\pi i x}, w \right) \right| \\ &= \frac{\pi N a_N^2}{1 + \delta \pi^2 N^2} \left| \left(e^{-\pi i x}, w \right) \right| \\ &= \frac{2\pi N a_N^2}{1 + \delta \pi^2 N^2}, \end{split}$$

the constant b(w) is much larger for very small δ and large N.

Obviously the main difficulties in the previous paragraph come from the high frequency terms. In order to remedy it, the restrain operator R_v is proposed in [13, 14] such that for

$$v(x) = \sum_{|n| \leq N} a_n e^{n\pi i x},$$

we have

$$R_{\nu}v(x) = \sum_{|n| \leq N} \left(1 - \left|\frac{n}{N}\right|^{\nu}\right) a_n e^{n\pi i x}.$$

In [14], the nonlinear term $U(x, t)(\partial/\partial x)$ U(x, t) is approximated by $J(u^k(x) + \sigma \tau u_t^k(x))$, $R_v u^k(x)$, which is successfully applied to the KDV equation. If we generalize this method to (1) directly, then we have

$$u_{t}^{k}(x) + \alpha \frac{\partial}{\partial x} u^{k}(x) + J(u^{k}(x) + \sigma \tau u_{t}^{k}(x), R_{v}u^{k}(x)) - \delta \frac{\partial^{2}}{\partial x^{2}} u_{t}^{k}(x)$$

$$= p_{c} f^{k}(x), \qquad x \in \vartheta, k \in S_{\tau},$$

$$u^{0}(x) = p_{c} U_{0}(x), \qquad x \in \vartheta.$$
(8)

When $\sigma = 0$, we can solve (10) explicitly. This is one of the advantages of the spectral method and pseudospectral methods. But the numerical result is unsatisfactory. In particular, the error is very big for $\alpha \neq 0$ and small δ (see Table VI, in Section III).

In this paper, we approximate the nonlinear term by $R_{\nu}J(R_{\nu}u^{k}(x) + \sigma\tau R_{\nu}u_{i}^{k}(x), u^{k}(x))$ and use the filtering technique for the linear term $\alpha(\partial/\partial x) u^{k}(x)$ and the right term $p_{c} f^{k}(x)$. Now we have the scheme

$$u_{t}^{k}(x) + \alpha R_{v} \frac{\partial}{\partial x} u^{k}(x) + R_{v} J(R_{v} u^{k}(x) + \sigma \tau R_{v} u_{t}^{k}(x), u^{k}(x))$$
$$- \delta \frac{\partial^{2}}{\partial x^{2}} u_{t}^{k}(x) = R_{v} p_{c} f^{k}(x), \quad x \in \vartheta. \quad k \in s_{\tau}, \qquad (9)$$
$$u^{0}(x) = p_{c} U_{0}(x), \quad x \in \vartheta.$$

Clearly,

$$\left(R_{\nu}\frac{\partial}{\partial x}v,v\right) = \left(\frac{\partial}{\partial x}R_{\nu}^{1/2}v,R_{\nu}^{1/2}v\right) = 0, \quad \forall v \in V_{N}.$$
(10)

From (5), for all $v, w, \varphi \in \dot{V}_N$,

$$(R_{v}J(R_{v}v,\varphi),w) + (R_{v}J(R_{v}w,\varphi),v) = 0.$$
(11)

If $\sigma = \frac{1}{2}$, then it follows from (9)–(11) that

$$\int_{0}^{2} u^{k}(x) dx = \int_{0}^{2} u^{0}(x) dx + \tau \sum_{\xi=0}^{k-1} \int_{0}^{2} R_{\nu} p_{c} f^{\xi}(x) dx,$$
$$\|u^{k}\|^{2} + \delta \|u^{k}\|_{1}^{2} = \|u^{0}\|^{2} + \delta \|u^{0}\|_{1}^{2} + \tau \sum_{\xi=0}^{k-1} (R_{\nu} p_{c} f^{\xi}, u^{\xi} + u^{\xi+1}),$$

which are the discrete analogy of (3) and (4).

III. NUMERICAL RESULTS

We test the effect of the restrain operator with the function

$$U(x, t) = A \exp(A \sin \pi x + 0.5t).$$

Put N = 8, $\tau = 0.01$, and let

$$E(t) = \max_{0 \le j \le 2N} \frac{|u(x_j, t) - U(x_j, t)|}{|U(x_j, t)|}.$$

All computations are carried out with the explicit scheme, i.e., $\sigma = 0$. The numerical results show that

(i) If the vibration of the genuine solution of (1) is small (e.g., A = 0.1) and $\alpha = 0$, then the effect of R_{ν} is not clear (see Table I). The accuracy of (9) is nearly the same as the spectral method of [11].

(ii) If δ is large (e.g., $\delta = 1$), then the operator R_{ν} is not important even though the vibration of a genuine solution is big (see Table II).

(iii) If the vibration of a genuine solution is big and δ small (e.g., A = 1, $\delta = 10^{-6}$, 10^{-4}), then the effect of R_{ν} is very clear, see Tables III and IV.

(iv) The smaller the parameter δ , the more important the linear term $\alpha(\partial U/\partial x)$. Thus the operator R_{ν} is very important even if the vibration of the genuine solution is not big. Tables I and IV show that the effect of R_{ν} is not clear for $\alpha = 0$, A = 0.1, but very clear for $\alpha = 2$, A = 0.1.

(v) The value of v in the restrain operator must be suitably chosen. If v is too large, the filtering technique is weakened. If v is too small, the approximation accuracy is lowered. The suitable value of v is between 5 and 10. But the best value of v is different in different cases. For instance, the best choice in Table IV is v = 3.

(vi) For the proof of convergence, the filtering technique for the linear term and right term is not important. But numerical experiment shows the importance, see Table V. In Table V, schemes $(9)_{a-c}$ are similar to (9), but the linear term and the right term are respectively approximated in the following

Scheme (9)a:
$$\alpha \frac{\partial}{\partial x} u^k(x),$$
 $R_v p_c f^k(x);$ Scheme (9)b: $\alpha R_v \frac{\partial}{\partial x} u^k(x),$ $p_c f^k(x);$ Scheme (9)c: $\alpha \frac{\partial}{\partial x} u^k(x),$ $p_c f^k(x).$

(vii) The numerical results of scheme (8) are much worse than that of scheme (9) and so the new filtering technique in this paper is better for the explicit scheme of a nonlinear problem. In Table VI, schemes (8) a-c are similar to scheme (8) with the linear term and the right term approximated by

Scheme (8)a: $\alpha R_v \frac{\partial}{\partial x} u^k(x), \quad p_c f^k(x);$ Scheme (8)b: $\alpha \frac{\partial}{\partial x} u^k(x), \quad R_v p_c f^k(x);$ Scheme (8)c: $\alpha \frac{\partial}{\partial x} u^k(x), \quad p_c f^k(x).$

TABLE I

Scheme (9), $\sigma = 0, \alpha = 0, A = 0.1$.

E(5.0)	v = 5	v = 10	v = 50	
$\delta = 10^{-6}$	0.00228242	0.00228203	0.00228193	
$\delta = 10^{-4}$	0.00228399	0.00228281	0.00228183	
$\delta = 10^{-2}$	0.00229817	0.00229622	0.00229651	
$\delta = 1$	0.00236857	0.00236493	0.00236557	

TABLE II

Scheme (9), $\sigma = 0$, $\delta = 1$, A = 1.

<i>E</i> (5.0)	v = 5	v = 10	v = 50
$\alpha = 2$	0.00144118	0.00118726	0.00115333
$\alpha = 0$	0.00231011	0.00215183	0.00216537

TABLE I	I	J
---------	---	---

Scheme (9), $\sigma = 0$, $\alpha = 0$, A = 1.

<i>E</i> (1.0)	v = 5	v = 10	v = 50
$\delta = 10^{-6}$ $\delta = 10^{-4}$	0.00329499	0.0355039	0.235286
	0.0283234	0.0124055	0.170894

Scheme (9), $\alpha = 2$, A = 0.1

<i>E</i> (2.0)	v = 3	v = 5	v = 10	v = 50
$\delta = 10^{-6}$ $\delta = 10^{-4}$	0.00156715	0.0186465	0.0243057	> 10
	0.00155333	0.00163698	0.0111043	0.702116

r	•	R	T	F	v
1	А	D	L	Е	¥

$$\sigma = 0, \ \alpha = 2, \ \nu = 10, \ \delta = 10^{-6}$$

	E(0.5), A = 1.	E(2.0), A = 0.1.
Scheme (9)	0.0454317	0.0243057
Scheme (9)a	0.0745903	7.42054
Scheme (9)b	0.564465	0.0275196
Scheme (9)c	2.08395	9.66944

	$\sigma = 0, \alpha = 2, A = 1, \nu = 10, \delta = 10^{-5}$				
	Scheme (8)	Scheme (8)a	Scheme (8)b	Scheme (8)c	
<i>E</i> (0.5)	0.136626	0.271632	7.86826	2.89614	

TABLE VI $\sigma = 0, \alpha = 2, A = 1, v = 10, \delta = 10^{-6}$

(viii) Another pseudospectral scheme is

-

$$u_{t}^{k}(x) + \alpha R_{v} \frac{\partial}{\partial x} u^{k}(x) + R_{v} J(R_{v} u^{k}(x) + \sigma \tau R_{v} u_{t}^{k}(x), R_{v} u^{k}(x))$$
$$- \delta \frac{\partial^{2}}{\partial x^{2}} u_{t}^{k}(x) = R_{v} p_{c} f^{k}(x).$$
(12)

This scheme seems more reasonable, but the numerical results show that the accuracy of (12) is nearly the same as scheme (9).

 $E(0.5), \alpha = 2$ $E(1.0), \alpha = 0$ $\delta = 10^{-6}$ $\delta = 10^{-4}$ $\delta = 10^{-6}$ $\delta = 10^{-4}$ Scheme (9) 0.0454317 0.0311898 0.0355039 0.0124055 Scheme (12) 0.0586913 0.0368031 0.0129624 0.0181770

TABLE VII $\sigma = 0, A = 1, v = 10$

TABLE	VIII
-------	------

 $\sigma = 0, A = 0.1, v = 10$

	<i>E</i> (2.0)), $\alpha = 2$	E(5.0)	$\alpha = 0$
	$\delta = 10^{-6}$	$\delta = 10^{-4}$	$\delta = 10^{-6}$	$\delta = 10^{-4}$
Scheme (9)	0.0243057	0.0111043	0.00228203	0.00228281
Scheme (12)	0.0244925	0.0124284	0.00228242	0.00228232

GUO AND CAO

IV. ERROR ESTIMATIONS

Assume that $u^0(x)$ and $f^k(x)$ have respectively the error $\tilde{u}^0(x)$ and $\tilde{f}^k(x)$. Then the error of $u^k(x)$, denoted by $\tilde{u}^k(x)$, satisfies

$$\tilde{u}_{t}^{k}(x) + \alpha R_{v} \frac{\partial}{\partial x} \tilde{u}^{k}(x) + R_{v} J(R_{v} \tilde{u}^{k}(x) + \sigma \tau R_{v} \tilde{u}_{t}^{k}(x), \tilde{u}^{k}(x) + u^{k}(x))$$

+ $R_{v} J(R_{v} u^{k}(x) + \sigma \tau R_{v} u_{t}^{k}(x), \tilde{u}^{k}(x)) - \delta \frac{\partial^{2}}{\partial x^{2}} \tilde{u}_{t}^{k}(x) = R_{v} p_{c} \tilde{f}^{k}(x).$ (13)

Let

$$L^{p}(\vartheta) = \left\{ v \mid ||v||_{L^{p}} = \left(\int_{0}^{2} |v(x)|^{p} dx \right)^{1/p} < \infty \right\}.$$

In particular, the inner product and the norm are (\cdot, \cdot) and $\|\cdot\|$, respectively, for p = 2. For any positive integer, let $\|v\|_{\beta} = \|\partial^{\beta} u/\partial x^{\beta}\|$,

$$H^{\beta}(\vartheta) = \left\{ v \left| \left\| v \right\|_{\beta} = \sum_{q=1}^{\beta} \left| v \right|_{q} < \infty \right\}.$$

For any positive constant β , $H^{\beta}(\vartheta)$ is the complex interpolation space between $H^{[\beta]}(\vartheta)$ and $H^{[\beta]+1}(\vartheta)$. Define

$$H^{\beta}_{p}(\vartheta) = \{ v \mid v \in H^{\beta}(\vartheta), v(x) = v(x+2) \}.$$

We denote by $C^{q}(0, T; H^{\beta}_{p}(\vartheta))$ the space of abstract functions with the norm

$$\|v\|_{C^{q}(0, T; H^{\beta}_{p}(\mathfrak{I}))} = \max_{0 \leq s \leq q} \max_{0 \leq t \leq T} \left\| \frac{\partial^{s} v(t)}{\partial t^{s}} \right\|_{\beta}$$

Similarly we define the space $H^q(0, T; H^{\beta}_p(\vartheta))$. Let $|||v|||_{\beta} = \max_{0 \le k\tau \le T} ||v^k||_{\beta}$. In particular, $|||v||| = \max_{0 \le k\tau \le T} ||v^k||$. Let p be a suitably small positive constant and

$$Q^{k}(v) = \|v^{k}\|^{2} + \delta \|v^{k}\|_{1}^{2} + p\tau^{2} \sum_{\xi=0}^{k-1} (\|v^{\xi}_{t}\|^{2} + \delta \|v^{\xi}_{t}\|_{1}^{2}),$$

$$\rho^{k}(v, w) = \|v\|^{2} + \delta \|v\|_{1}^{2} + \tau \sum_{\xi=0}^{k-1} \|w^{k}\|^{2}.$$

THEOREM 1. (i) If $\sigma \leq \frac{1}{2}$, then there exist positive constants b_1, b_2, b_3 depending only on $|||u||_{3/2+\gamma}$ ($\gamma > 0$) and δ such that for all $k\tau \leq T$, $\rho^k(\tilde{u}^0, \tilde{f}) \leq b_1/\tau N$,

$$Q^k(\tilde{u}) \leq b_2 \rho^k(\tilde{u}^0, \tilde{f}) e^{b_3 k \tau}$$

(ii) If $\sigma > \frac{1}{2}$ then there exist positive constants b_4 , b_5 depending on $|||u|||_{3/2+\gamma}$ and δ such that for all k and $\rho^k(\tilde{u}^0, \tilde{f})$,

$$Q^k(\tilde{u}) \leq b_4 \rho^k(\tilde{u}^0, \tilde{f}) e^{b_5 k \tau}.$$

(iii) If $\sigma \leq \frac{1}{2}$, $\tau = O(N^{-2})$, then there exist positive constants b_6 , b_7 , b_8 depending only on $|||u|||_{3/2+\gamma}$ such that for all $k\tau \leq T$, $\rho^k(\tilde{u}^0, \tilde{f}) \leq b_6 N^{-1}$,

$$Q^k(\tilde{u}) \leq b_7 \rho^k(\tilde{u}^0, \tilde{f}) e^{b_8 k \tau}$$

(iv) If $\sigma > \frac{1}{2}$, $\tau = O(N^{-2})$, then there exist positive constants b_9 , b_{10} depending only on $|||u|||_{3/2+\gamma}$ such that for all k and $\rho^k(\tilde{u}^0, \tilde{f})$,

$$Q^k(\tilde{u}) \leqslant b_9 \rho^k(\tilde{u}^0, \tilde{f}) e^{b_{10}k\tau}.$$

Next consider the convergence. Let U(x, t) and $u^k(x)$ be the solutions of (1) and (9), respectively.

THEOREM 2. Assume that

(i) $U \in H^2(0, T; H^1_p(\vartheta)) \cap C(0, T; H^\beta_p(\vartheta)), U_0 \in H^\beta_p(\vartheta), f \in C(0, T; H^{\beta-1}_p(\vartheta));$ (ii) $v \ge \beta \ge 2$,

then for all $k\tau \leq T$,

$$Q^{k}(u-U) \leq b_{11}e^{b_{12}k\tau}(\tau^{2}+N^{2-2\beta}),$$

where b_{11} and b_{12} are positive constants depending on $||U_6||_{\beta}$, $||f||_{c(0, T; H^{\beta-1}(\beta))}$, $||U||_{H^2(0, T; H^1(\beta))}$, $||U||_{c(0, T; H^{\beta}(\beta))}$, and δ . If $\tau = O(N^{-2})$, then b_{11} and b_{12} are independent of δ .

V. SOME LEMMAS

For the proof of the theorems, we need the following lemmas.

LEMMA 1 [12]. If $0 \le \mu \le \beta$, $v \in H_p^{\beta}(\vartheta)$, then

$$\|P_N v - v\|_{\mu} \leq c N^{\mu - \beta} \|v\|_{\beta}.$$

If $\beta > \frac{1}{2}$, in addition, then

$$\|\boldsymbol{P}_{c}\boldsymbol{v}-\boldsymbol{v}\|_{\mu} \leq cN^{\mu-\beta}\|\boldsymbol{v}\|_{\beta}.$$

LEMMA 2 [12]. If $0 \le \mu \le \beta$, $v \in V_N$, then

$$\|v\|_{\beta} \leq c N^{\beta-\mu} \|v\|_{\mu}.$$

LEMMA 3 [12]. If $v, w \in V_N$, then

$$|vw||^2 \leq (2N+1) ||v||^2 ||w||^2$$
.

LEMMA 4 [14]. If $0 \le \mu \le \beta \le v$, $v \in V_N$, then

$$\|R_{v}v-v\|_{\mu} \leq cN^{\mu-\beta} |v|_{\beta}.$$

For simplicity, we introduce the circle convolution operator "*" such that if

$$v(x) = \sum_{|n| \leq N} a_n e^{n\pi i x}, \qquad w(x) = \sum_{|n| \leq N} b_n e^{n\pi i x},$$

then

$$v * w(x) = \sum_{|n| \leq N} \sum_{|l| \leq N} a_l b_{n-l} e^{n\pi i x}$$

where $a_{n+2N+1} = a_n$, $b_{n+2N+1} = b_n$. From [14], for $v, w \in V_N$, $\varphi \in \dot{V}_N$, we have

$$P_c(v(x) w(x)) = v * w(x),$$
 (14)

$$(v * \varphi, w) = (v, \varphi * \omega). \tag{15}$$

LEMMA 5 [14]. If $v \in V_N$, $w \in H_p^{3/2+\gamma}(\vartheta)$, $\gamma > 0$, then $\left| \left(R_v v * \frac{\partial v}{\partial x}, w \right) \right| \leq c_\gamma v \|w\|_{3/2+\gamma} \|v\|^2,$ $\left| \left(v * R_v \frac{\partial v}{\partial x}, w \right) \right| \leq c_\gamma v \|w\|_{3/2+\gamma} \|v\|^2,$

where c_{γ} is a positive constant depending only on γ .

LEMMA 6. If $v, w \in V_N, \gamma > 0$, then

$$|(R_{v}J(R_{v}w, v), v)| \leq c_{y} ||w||_{3/2+y} ||v||^{2}.$$

Proof. We have

$$J(R_{v}w, v) = \frac{1}{3}v * R_{v}\frac{\partial w}{\partial x} + \frac{1}{3}\frac{\partial}{\partial x}(v * R_{v}w),$$

$$|(R_{v}J(R_{v}w, v), v)| = |(J(R_{v}\omega, v), R_{v}v)|$$

$$\leq \frac{1}{3}\left|\left(v * R_{v}\frac{\partial w}{\partial x}, R_{v}v\right)\right| + \frac{1}{3}\left|\left(\frac{\partial}{\partial x}(v * R_{v}w), R_{v}v\right)\right|.$$

From (15) and Lemma 5

$$\left| \left(v * R_{v} \frac{\partial w}{\partial x}, R_{v} v \right) \right| \leq c \left\| \frac{\partial w}{\partial x} \right\|_{L^{\infty}(9)} \|v\|^{2} \leq c_{y} \|w\|_{3/2 + y} \|v\|^{2},$$
$$\left| \left(\frac{\partial}{\partial x} \left(v * R_{v} w \right), R_{v} v \right) \right| = \left| \left(v * R_{v} \frac{\partial v}{\partial x}, R_{v} w \right) \right| \leq c_{y} \|w\|_{3/2 + y} \|v\|^{2}.$$

This completes the proof.

LEMMA 7 [15]. If the following conditions are fulfilled

(i) Z^k is nonnegative function of k, M_1, M_2, M_3 , and ρ are nonnegative constants,

- (ii) α_1, α_2 , and α_3 are constants,
- (iii) if $Y \leq M_1 N^{\alpha_1}$, then $\eta(Y) \leq 0$,
- (iv) for all $k \ge 1$,

$$Z^{k} \leq \rho + \tau \sum_{\xi=0}^{k-1} \left[M_2 Z^{\xi} + M_3 (Z^{\xi})^{\alpha_2} N^{\alpha_3} + \eta (Z^{\xi}) \right],$$

(v) $Z^0 \leq \rho \leq e^{-(M_2 + M_3)T} \min(N^{-\alpha_3/\alpha_2}, M_1 N^{\alpha_1}),$

then for all $k\tau \leq T$,

$$Z^k \leqslant \rho \ e^{(M_2 + M_3)k\tau}. \tag{16}$$

In particular, (16) holds for all k and ρ provided $M_3 = 0$ and $\eta(Y) \leq 0$.

VI. THE PROOF OF THEOREM 1

By taking the inner product of (13) with $2\tilde{u}^k(x)$, we have from (10) and (11) that

$$(\|\tilde{u}^{k}\|^{2} + \delta |\tilde{u}^{k}|_{1}^{2})_{t} - \tau(\|\tilde{u}^{k}_{t}\|^{2} + \delta |\tilde{u}^{k}_{t}|_{1}^{2}) - 2\sigma\tau(R_{v}J(R_{v}\tilde{u}^{k}, u^{k}, u^{k} + \tilde{u}^{k}), \tilde{u}^{k}) + 2(R_{v}J(R_{v}u^{k} + \sigma\tau R_{v}u^{k}_{t}, \tilde{u}^{k}), \tilde{u}^{k})$$

$$= 2(R_{v}p_{c}\tilde{f}^{k}, \tilde{u}^{k}).$$
(17)

Let *m* and ε be positive constants. Taking the inner product of (13) with $m\tau \tilde{u}_t^k$, it follows from (11) that

$$m\tau(\|\tilde{u}_{t}^{k}\|^{2} + \delta |\tilde{u}_{t}^{k}|_{1}^{2}) + \alpha m\tau\left(R_{v}\frac{\partial}{\partial x}\tilde{u}^{k}, \tilde{u}_{t}^{k}\right) + m\tau(R_{v}J(R_{v}\tilde{u}^{k}, u^{k} + \tilde{u}^{k}),$$
$$\tilde{u}_{t}^{k}) + m\tau(R_{v}J(R_{v}u^{k} + \sigma\tau R_{v}u_{t}^{k}, \tilde{u}^{k}), \tilde{u}_{t}^{k}) = m\tau(R_{v}p_{c}\tilde{f}^{k}, \tilde{u}_{t}^{k}).$$
(18)

The combination of (17) with (18) leads to

$$(\|\tilde{u}^{k}\|^{2} + \delta \|\tilde{u}^{k}\|_{1}^{2})_{t} + \tau(m - 1 - \varepsilon)(\|\tilde{u}^{k}_{t}\|^{2} + \delta \|\tilde{u}^{k}_{t}\|_{1}^{2}) + \sum_{\ell=1}^{4} F_{\ell}^{k}$$

$$\leq \|\tilde{u}^{k}\|^{2} + \left(1 + \frac{\tau m^{2}}{4\varepsilon}\right) \|p_{c}\tilde{f}^{k}\|^{2}, \qquad (19)$$

where

$$F_1^k = \tau(m - 2\sigma)(R_v J(R_v \tilde{u}^k, u^k + \tilde{u}^k), \tilde{u}_t^k),$$

$$F_2^k = (R_v J(R_v u^k + \sigma \tau R_v u_t^k, \tilde{u}^k), \tilde{u}^k),$$

$$F_3^k = m\tau(R_v J(R_v u^k + \sigma \tau R_v u_t^k, \tilde{u}^k), \tilde{u}_t^k),$$

$$F_4^k = \alpha m\tau \left(R_v \frac{\partial}{\partial x} \tilde{u}^k, \tilde{u}_t^k\right).$$

Now we estimate $|F_l^k|$. From Lemmas 1-4, we have

$$\begin{split} |F_{1}^{k}| &\leq \varepsilon\tau \|\tilde{u}_{1}^{k}\|^{2} + \frac{c\tau(m-2\sigma)^{2}}{\varepsilon} \left(\left\| \frac{\partial u^{k}}{\partial x} \right\|_{L^{\infty}(\vartheta)}^{2} \|\tilde{u}^{k}\|^{2} + \|u^{k}\|_{L^{\infty}(\vartheta)} \|\tilde{u}^{k}\|_{1}^{2} \right. \\ &+ N \|\tilde{u}^{k}\|^{2} |\tilde{u}^{k}|_{1}^{2} \right) \\ &\leq \varepsilon\tau \|\tilde{u}_{1}^{k}\|^{2} + \frac{c\tau(m-2\sigma)^{2}}{\varepsilon} \left(\|u^{k}\|_{3/2+\gamma} (\|\tilde{u}^{k}\|^{2} + |\tilde{u}^{k}|_{1}^{2}) + N \|\tilde{u}^{k}\|^{2} |\tilde{u}^{k}|_{1}^{2} \right). \end{split}$$

Lemma 6 leads to

$$|F_2^k| \leq c |||u|||_{3/2+\gamma} ||\tilde{u}^k||^2.$$

We have also

$$\begin{split} |F_{3}^{k}| &\leq \varepsilon \tau \|\tilde{u}_{t}^{k}\|^{2} + \frac{c\tau m^{2}}{\varepsilon} \|\|u\|\|_{3/2 + \gamma}^{2} (\|\tilde{u}^{k}\|^{2} + |\tilde{u}^{k}|_{1}^{2}), \\ |F_{4}^{k}| &\leq \varepsilon \tau \|\tilde{u}_{t}^{k}\|^{2} + \frac{\tau \alpha^{2} m^{2}}{\varepsilon} \|\tilde{u}^{k}\|_{1}^{2}. \end{split}$$

Substituting the above estimations into (19), we get

$$(\|\tilde{u}^{k}\|^{2} + \delta |\tilde{u}^{k}|_{1}^{2})_{t} + \tau (m - 1 - 4\varepsilon)(\|\tilde{u}^{k}_{t}\|^{2} + \delta |\tilde{u}^{k}_{t}|_{1}^{2}) \leq R^{k},$$
(20)

where

$$R^{k} = c \left\{ \left[1 + \frac{\tau m^{2}}{\varepsilon} + \left(1 + \frac{\tau (m - 2\sigma)^{2}}{\varepsilon} \right) \|\|u\|\|_{3/2 + \gamma}^{2} \right] \|\tilde{u}^{k}\|^{2} + \frac{\tau (m - 2\sigma)^{2} + \tau m^{2}}{\varepsilon} \\ \times (\|\|u\|\|_{3/2 + \gamma}^{2} + 1) \|\tilde{u}^{k}\|_{1}^{2} + \frac{\tau N(m - 2\sigma)^{2}}{\varepsilon} \|\tilde{u}^{k}\|^{2} \|\tilde{u}^{k}\|_{1}^{2} + \left(1 + \frac{\tau m^{2}}{4\varepsilon} \right) \|\tilde{f}^{k}\|^{2} \right\}.$$

Let p be suitably small. Take $\varepsilon = \frac{1}{4}(m-1-p) > 0$ and

$$\begin{split} m &= 2\sigma, \qquad H_{\sigma} = 0, \qquad \text{for} \quad \sigma > \frac{1}{2}, \\ m &= 2, \qquad H_{\sigma} = 1, \qquad \text{for} \quad \sigma \leqslant \frac{1}{2}. \end{split}$$

Then

$$(\|\tilde{u}^{k}\|^{2} + \delta |\tilde{u}^{k}|_{1}^{2})_{t} + p\tau(\|\tilde{u}^{k}_{t}\|^{2} + \delta |\tilde{u}^{k}_{t}|_{1}^{2}) \leq \mathbb{R}^{k},$$
(21)

from which

$$Q^{k}(\tilde{u}) \leq c\rho^{k}(\tilde{u}^{0},\tilde{f}) + c\tau \sum_{\xi=0}^{k-1} \left[(1 + |||u|||_{3/2+\gamma}^{2}) (||\tilde{u}^{\xi}||^{2} + \delta |\tilde{u}^{\xi}|_{1}^{2}) + \tau NH_{\sigma} ||\tilde{u}^{\xi}||^{2} |\tilde{u}^{\xi}|_{1}^{2} \right].$$

Since Lemma 2 and

$$N \|\tilde{u}^{k}\|^{2} \|\tilde{u}^{k}\|_{1}^{2} \leq \delta N \|\tilde{u}^{k}\|_{1}^{4} + \frac{N}{4\delta} \|\tilde{u}^{k}\|^{4}$$

or

$$N \|\tilde{u}^k\|^2 \|\tilde{u}^k\|_1^2 \leq c N^3 \|\tilde{u}^k\|^4.$$

Lemma 7 completes the proof.

VII. THE PROOF OF THEOREM 2

Let $W(x, t) = p_N U(x, t)$. From (1), we have

$$\frac{\partial W(x,t)}{\partial t} + \alpha \frac{\partial W(x,t)}{\partial x} + p_N \left(U(x,t) \frac{\partial U(x,t)}{\partial x} \right) - \delta \frac{\partial^3 W(x,t)}{\partial t \partial x^2} = p_N f(x,t)$$

and so

$$W_{t}^{k}(x) + \alpha R_{v} \frac{\partial W^{k}(x)}{\partial x} + R_{v} J(R_{v} W^{k}(x) + \sigma \tau R_{v} W_{t}^{k}(x), W^{k}(x))$$
$$-\delta \frac{\partial^{2} W_{t}^{k}(x)}{\partial x^{2}} = R_{v} p_{c} f^{k}(x) + \sum_{l=1}^{5} g_{l}^{k}(x), \qquad (22)$$

where

$$g_1^k(x) = W_1^k(x) - \frac{\partial W(x, k\tau)}{\partial t},$$

$$g_2^k(x) = \alpha R_v \frac{\partial W^k(x)}{\partial x} - \alpha \frac{\partial W^k(x)}{\partial x},$$

$$g_3^k(x) = R_v J(R_v W^k(x) + \sigma \tau R_v W_1^k(x), W^k(x)) - p_N \left(U(x, k\tau) \frac{\partial U(x, k\tau)}{\partial x} \right),$$

$$g_4^k(x) = \delta \frac{\partial^3 W^k(x)}{\partial t \partial x^2} - \delta \frac{\partial^2 W_1^k(x)}{\partial x^2},$$

$$g_5^k(x) = p_N f^k(x) - R_v p_c f^k(x).$$

Let $e^k(x) = u^k(x) - W^k(x)$. From (9) and (22), it follows that

$$e_t^k(x) + \alpha R_v \frac{\partial e^k(x)}{\partial x} + R_v J(R_v e^k(x) + \sigma \tau R_v e_t^k(x), W^k(x) + e^k(x)) + R_v J(R_v W^k(x) + \sigma \tau R_v W_t^k(x), e^k(x)) - \delta \frac{\partial^2 e_t^k(x)}{\partial x^2} = -\sum_{l=1}^5 g_l^k(x).$$

Integration by parts leads to

$$\frac{\partial W(x,k\tau)}{\partial t} - W_t^k(x) = -\frac{1}{\tau} \int_{k\tau}^{(k+1)\tau} (k\tau + \tau - \eta) \frac{\partial^2 W(x,\eta)}{\partial \eta^2} d\eta$$

and thus

$$\tau \sum_{\xi=0}^{k-1} \|g_1^{\xi}\|^2 \leq c\tau^2 \|U\|_{H^2(0, T; L^2(\vartheta))}^2.$$

It is easy to prove that

$$\tau \sum_{\xi=0}^{k-1} \|g_2^{\xi}\|^2 \leq c N^{2-2\beta} \|U\|_{c(0, T; H^{\beta}(\mathcal{G}))}^2.$$

By an argument similar to that of [14],

$$\tau \sum_{\xi=0}^{k-1} \|g_{3}^{\xi}\|^{2} \leq c N^{2-2\beta} \|U\|_{c(0,T; H^{\beta}(\beta))}^{4}.$$

We have

$$\begin{aligned} \left| \tau \sum_{\xi=0}^{k-1} \left(g_{4}^{\xi}, e^{\xi} \right) \right| &\leq \delta \tau \sum_{\xi=0}^{k-1} \left| \left(\frac{\partial W_{t}^{\xi}}{\partial x} - \frac{\partial^{2} W(\xi \tau)}{\partial t \partial x}, \frac{\partial e^{\xi}}{\partial x} \right) \right| \\ &\leq \delta \tau \sum_{\xi=0}^{k-1} \left(\left| e^{\xi} \right|_{1}^{2} + \left\| \frac{\partial W_{t}^{\xi}}{\partial x} - \frac{\partial^{2} W(\xi \tau)}{\partial t \partial x} \right\| \right) \\ &\leq \delta \tau \sum_{\xi=0}^{k-1} \left| e^{\xi} \right|_{1}^{2} + c \delta \tau^{2} \left\| U \right\|_{H^{2}(0, T; H^{1}(\theta))}^{2}, \\ \left| m \tau^{2} \sum_{\xi=0}^{k-1} \left(g_{4}^{\xi}, e^{\xi}_{t} \right) \right| &\leq \varepsilon \, \delta \tau^{2} \sum_{\xi=0}^{k-1} \left| e^{\xi} \right|_{1}^{2} + c \delta \tau^{3} \left\| U \right\|_{H^{2}(0, T; H^{1}(\theta))}^{2}, \end{aligned}$$

We have also

$$\tau \sum_{\xi=0}^{k-1} \|g_{\xi}^{\xi}\|^{2} \leq c N^{2-2\beta} \|f\|_{c(0,T;H^{\beta-1}(\vartheta))}^{2},$$
$$\|e^{0}\|_{1}^{2} \leq c N^{2-2\beta} \|U_{0}\|_{\beta}^{2}.$$

The previous estimations give

$$\begin{split} \|e^{0}\|_{1}^{2} + \tau \sum_{\xi=0}^{k-1} \left(\sum_{l=1,2,3,5} \|g_{l}^{\xi}\|^{2} + |(g_{4}^{\xi}, e^{\xi} + m\tau e_{l}^{\xi})| \right) \\ &\leqslant \delta \tau \sum_{\xi=0}^{k-1} \left(|e^{\xi}|_{1}^{2} + \varepsilon \tau |e_{l}^{\xi}|_{1}^{2} \right) + c\tau^{2} \|U\|_{H^{2}(0,T; L^{2}(9))}^{2} + c\tau^{2} \|U\|_{H^{2}(0,T; H^{1}(9))}^{2} \\ &+ cN^{2-2\beta} \|U\|_{c(0,T; H^{\beta}(9))}^{4} + cN^{2-2\beta} \|U_{0}\|_{\beta}^{2} + cN^{2-2\beta} \|f\|_{c(0,T; H^{\beta-1}(9))}^{2}. \end{split}$$

Finally we finish the proof with an argument similar to that in Section VI and with the use of the triangular inequality.

References

- 1. D. H. PEREGRINE, J. Fluid Mech. 25, 321 (1966).
- 2. KH. O. ABDULLOEV, I. L. BOGOLUBSKY, AND V. T. MAKHANKOV, Phys. Lett. A 56, 427 (1976).
- 3. P. J. OLVER, Math. Proc. Cambridge Philos. Soc. 85, 143 (1979).
- 4. J. C. EILBECK AND G. R. MCGUIRE, J. Comput. Phys. 19, 43 (1975).
- 5. J. C. EILBECK AND G. R. MCGUIRE, J. Comput. Phys. 23, 63 (1977).
- 6. M. E. ALEXANDER AND J. L. MORRIS, J. Comput. Phys. 30, 428 (1979).
- 7. D. N. ARNOLD, J. DOUGLAS, AND V. THOMEE, Math. Comput. 36, 53 (1981).
- 8. H. M. WU AND B. Y. GUO, Math. Numer. Sin. 5, 90 (1983).

- 9. D. GOTTLIEB AND S. A. ORZSAG, "Numerical Analysis of Spectral Methods," CBMS-NSF, Regional Conference Series in Applied Mathematics, Vol. 26 (1977).
- 10. J. E. PASCIAK, SIAM J. Numer. Anal. 19, 142 (1982).
- 11. B. Y. GUO AND V. S. MANORANJAN, IMA-JNA 5, 307 (1985).
- 12. H. O. KREISS AND J. OLIGER, SIAM J. Numer. Anal. 16, 421 (1979).
- 13. P. Y. KUO, J. Comput. Math. 1, 353 (1983).
- 14. H. P. MA AND B. Y. GUO, J. Comput. Phys. 65, 120 (1986).
- 15. B.-Y. Guo, Acta Math. Sin. 17, 242 (1974).
- 16. B.-Y. Guo, Sci. Sin. A 25, 702 (1982).
- 17. D. F. GRIFFITHS, Bull. Inst. Math. Appl. 12, 210 (1982).